

Phi meson propagation in a hot hadronic gas

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Introduction

Quark-Gluon Plasma



Enhanced production of strange quarks



Enhancement of strange particles (K , ϕ , Λ , Ω)

Introduction

ϕ is a nice probe: \leftarrow A. Shor, P.R.L. 54 (1985) 1122

- ϕ production in pp and πp is **OZI suppressed**
- Coalescence: $s\bar{s} \rightarrow \phi$
- ϕ is not masked behind other resonances
- it can be detected using both kaon pairs ($K^+ K^-$) and dileptons ($e^+ e^-$, $\mu^+ \mu^-$)

Introduction

- Widely accepted that ϕ 's have a large mean free path in a hot hadronic matter

For example (K. Haglin, N.P.A 584 (1993) 719):

$$\lambda (T = 200 \text{ MeV}) = 4.4 \text{ fm}$$

$$\lambda (T = 150 \text{ MeV}) = 14 \text{ fm}$$



Phenomenological Lagrangians with couplings extracted from observed partial decay rates



Many mechanisms that include vertices allowed by the symmetries of QCD, like $\phi K^* K$ or $\rho K^* K^*$, are not taken into account.

Hidden Local Symmetry Lagrangian

- **Vector** mesons ρ , ω , K^* , ϕ are the **gauge bosons** of the hidden local $U(3)_V$ symmetry.

M. Bando, T. Kugo, K. Yamawaki, P.Rep. 164 (1988) 217

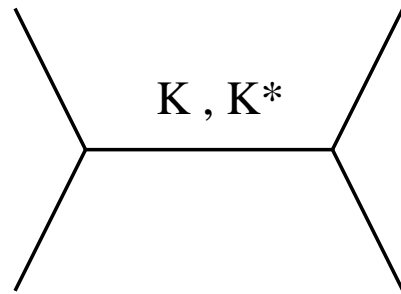
- Spontaneous symmetry breaking $\Rightarrow m_V$

$$m_V^2 = 2g^2 f_\pi^2, \quad g^2 \leftarrow VPP \text{ coupling}$$

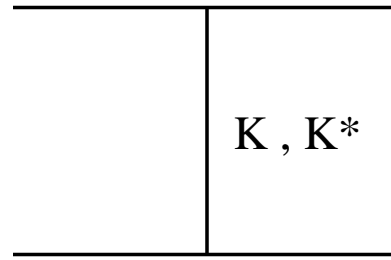
- Flavor symmetry breaking \Rightarrow different couplings
- Anomalous part of the Lagrangian $\Rightarrow VVP$ vertex
- Kinetic term for **vector** mesons \Rightarrow
 VVV and $VVVV$ vertices:

Binary reactions

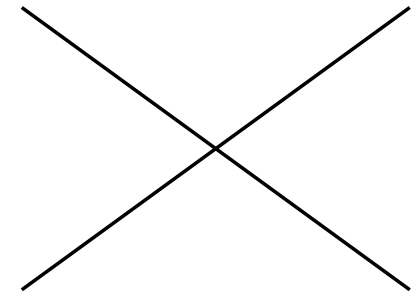
- Relevant vertices: VPP , VVP , VVV , $VVVV$



(s)



(t)



(c)

- No direct couplings of ϕ with π , ρ or ω (OZI rule)

Binary reactions

No.	Reaction	Channels
1.1	$\phi + \pi \rightarrow K + K$	$t(K^*)$
1.2	$\phi + \pi \rightarrow K + K^*$	$t(K, K^*)$
1.3	$\phi + \pi \rightarrow K^* + K^*$	$t(K, K^*)$

Binary reactions

No.	Reaction	Channels
2.1	$\phi + K \rightarrow \pi + K$	$\mathbf{s}, \mathbf{t}(K^*)$
2.2	$\phi + K \rightarrow \rho + K$	$\mathbf{s}, \mathbf{t}(K, K^*)$
2.3	$\phi + K \rightarrow \omega + K$	$\mathbf{s}, \mathbf{t}(K, K^*)$
2.4	$\phi + K \rightarrow \phi + K$	$\mathbf{s}, \mathbf{t}(K, K^*)$
2.5	$\phi + K \rightarrow \pi + K^*$	$\mathbf{s}(K, K^*), \mathbf{t}(K^*)$
2.6	$\phi + K \rightarrow \rho + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*)$
2.7	$\phi + K \rightarrow \omega + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*)$
2.8	$\phi + K \rightarrow \phi + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*)$

Binary reactions

No.	Reaction	Channels
3.1	$\phi + \rho \rightarrow K + K$	$t(K, K^*)$
3.2	$\phi + \rho \rightarrow K + K^*$	$t(K, K^*)$
3.3	$\phi + \rho \rightarrow K^* + K^*$	$t(K, K^*), c$

Binary reactions

No.	Reaction	Channels
4.1	$\phi + \omega \rightarrow K + K$	$t(K, K^*)$
4.2	$\phi + \omega \rightarrow K + K^*$	$t(K, K^*)$
4.3	$\phi + \omega \rightarrow K^* + K^*$	$t(K, K^*), c$

Binary reactions

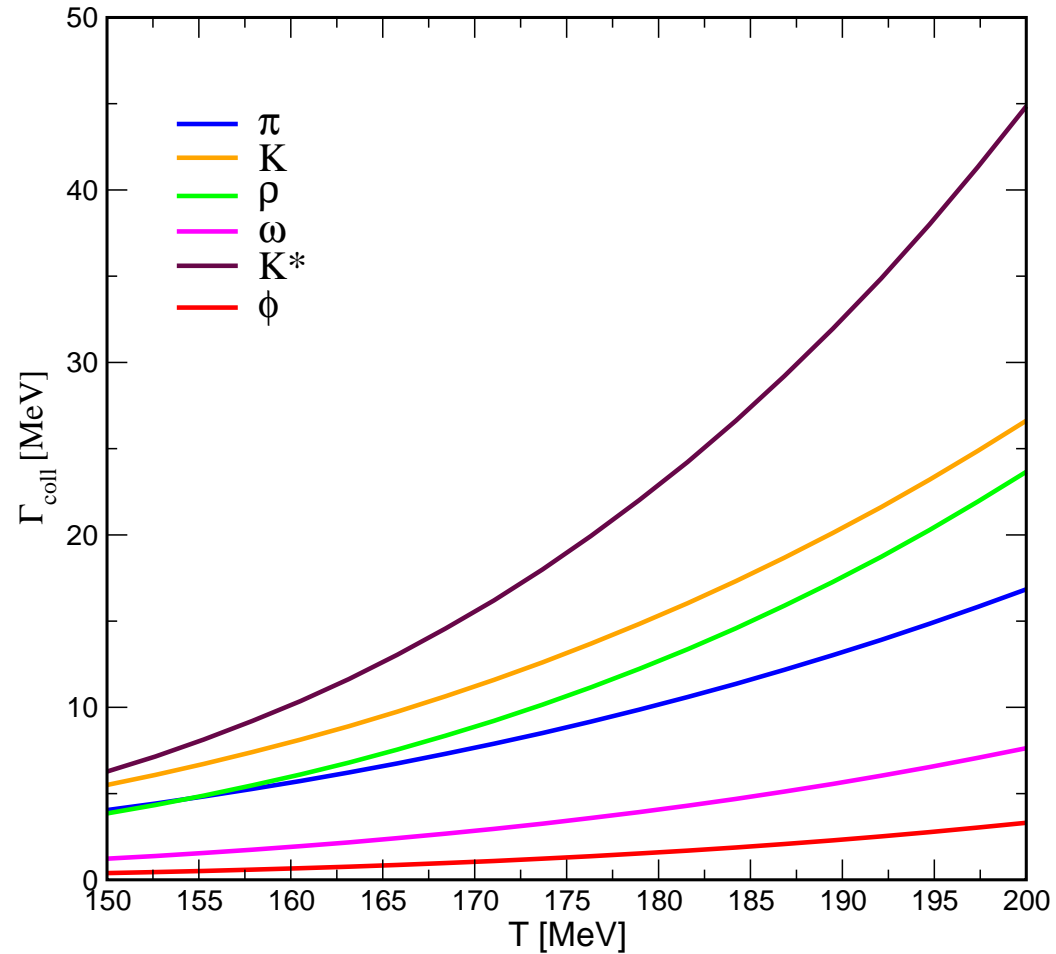
No.	Reaction	Channels
5.1	$\phi + K^* \rightarrow \pi + K$	$\mathbf{s}(K^*), \mathbf{t}(K, K^*)$
5.2	$\phi + K^* \rightarrow \rho + K$	$\mathbf{s}, \mathbf{t}(K, K^*)$
5.3	$\phi + K^* \rightarrow \omega + K$	$\mathbf{s}, \mathbf{t}(K, K^*)$
5.4	$\phi + K^* \rightarrow \phi + K$	$\mathbf{s}, \mathbf{t}(K, K^*)$
2.5	$\phi + K^* \rightarrow \pi + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*)$
5.6	$\phi + K^* \rightarrow \rho + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*), \mathbf{c}$
5.7	$\phi + K^* \rightarrow \omega + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*), \mathbf{c}$
5.8	$\phi + K^* \rightarrow \phi + K^*$	$\mathbf{s}, \mathbf{t}(K, K^*), \mathbf{c}$

Binary reactions

No.	Reaction	Channels
6.1	$\phi + \phi \rightarrow K + K$	$t(K, K^*)$
6.2	$\phi + \phi \rightarrow K + K^*$	$t(K, K^*)$
6.3	$\phi + \phi \rightarrow K^* + K^*$	$t(K, K^*), c$

Collision rates

$$\Gamma_{coll}^{(a)} = n_a \langle \sigma_{\phi a} v_{rel} \rangle \quad a = \pi, K, \rho, \omega, K^*, \phi$$



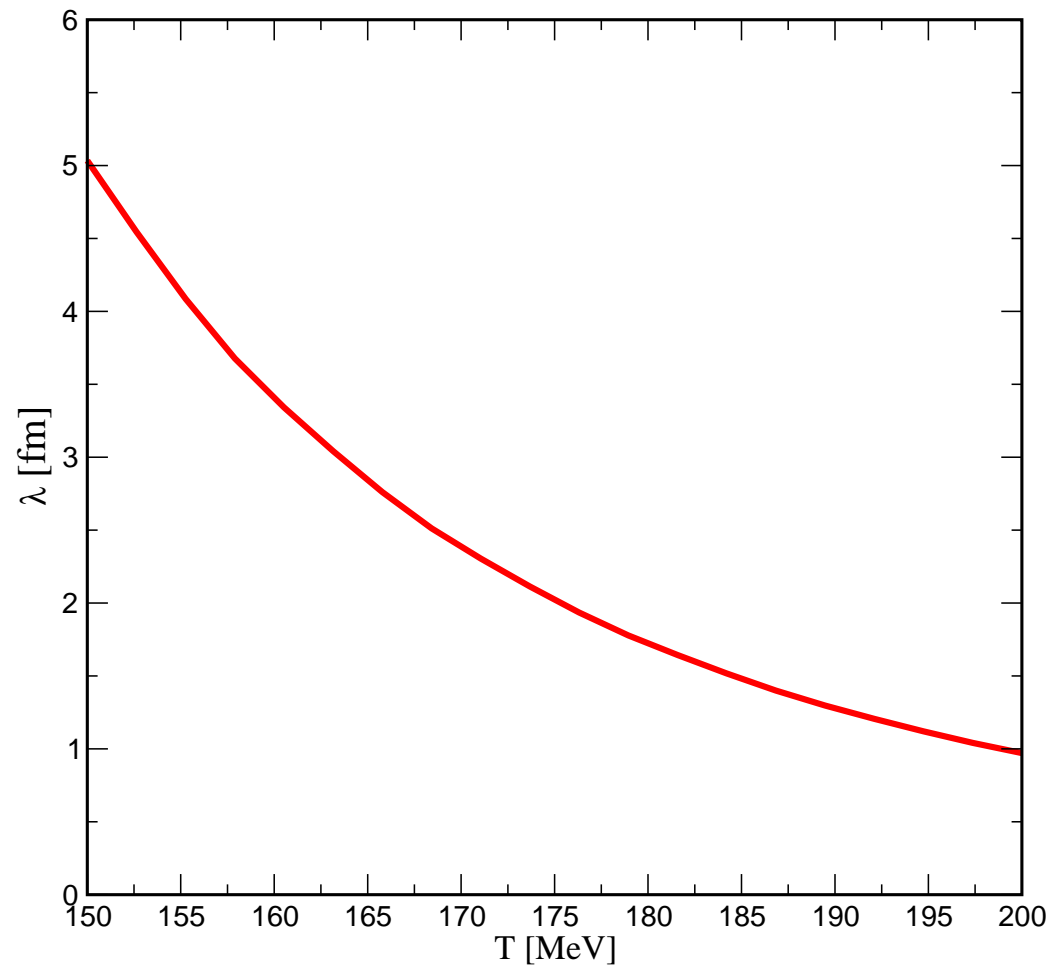
Collision rates

- Why $\Gamma_{coll}^{(K^*)} > \Gamma_{coll}^{(K)}$?

$$\frac{n_{K^*}}{n_K} = \frac{3m_{K^*}^2 K_2(m_{K^*}/T)}{m_K^2 K_2(m_K/T)} \Big|_{T=200 \text{ MeV}} = 0.77 \leftarrow \text{not so small}$$

- Large contributions to $\sigma_{\phi K^*}$ from:
 $\phi + K^* \rightarrow \rho + K^*$
 $\phi + K^* \rightarrow \omega + K^*$
 $\phi + K^* \rightarrow \phi + K^*$
- Inelastic reactions account for **> 80 %** of Γ_{coll}

Mean free path



Very short mean free path !

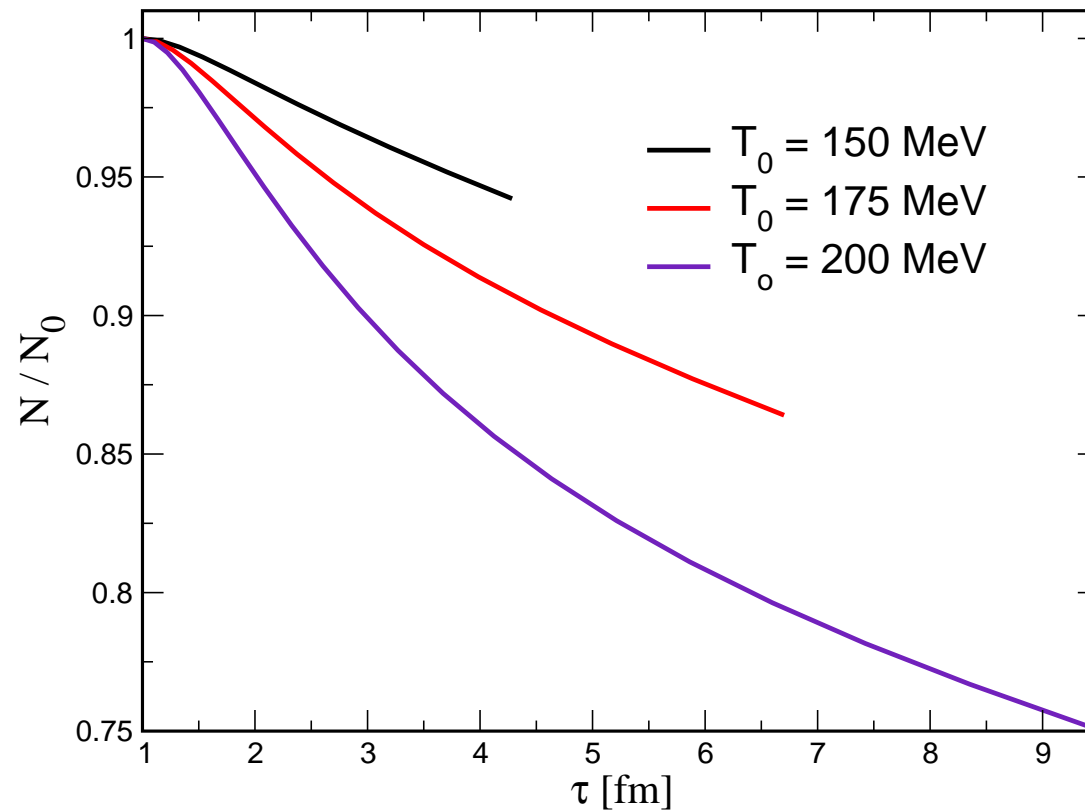
Time evolution of ϕ number in an expanding hadronic fireball

Assumptions:

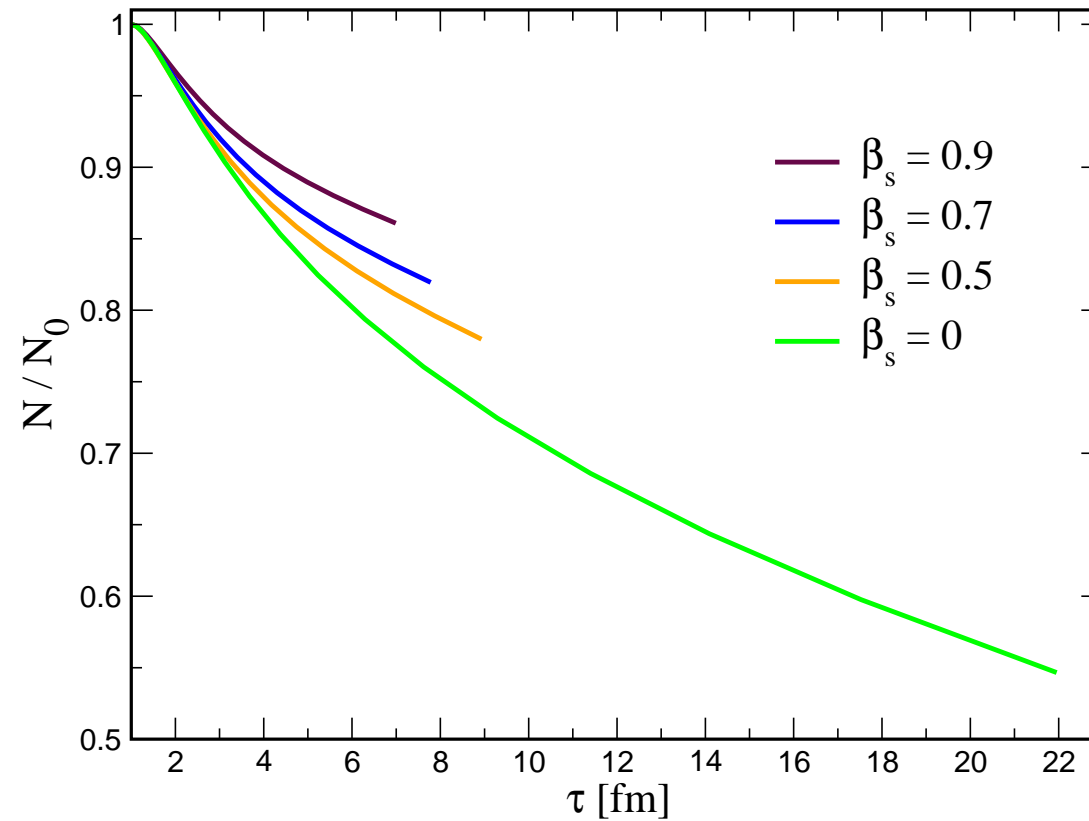
- An **ideal gas** of π , K , ρ , ω , K^*
- Kinetic equilibrium for all species including ϕ 's
- Chemical equilibrium for all species **except** ϕ 's
- Cylindrical symmetry and boost invariance
- Longitudinal and radial transverse expansion
- Entropy conservation
- Rate equation:

$$\partial_\mu (n_\phi u^\mu) = -\Gamma (n_\phi - n_\phi^{eq}), \quad \Gamma = \Gamma_{coll}(T) + \Gamma_{dec}(T)$$

Different hadronization temperatures



Different flow velocities



Conclusions

- ϕ meson interaction with a hot hadronic gas composed of π , K , ρ , ω , K^* , ϕ has been studied using the Hidden Local Symmetry model.
- Large ϕ collision rates with K , ρ and specially K^* .
- Short mean free path:
 $1 \leq \lambda \leq 5$ fm at $200 \geq T \geq 150$ MeV.
- High collision rates cause a reduction of the ϕ number.
- $N/N_0 \sim 5 - 45$ % depending on hadronization and freezeout temperatures, flow velocity and chemical potentials.